

## MSci EXAMINATION

### PHY-966(4261) Electromagnetic Theory

Time Allowed: 2 hours 30 minutes

Date:

Time:

Instructions: Answer **THREE QUESTIONS** only. Each question carries 20 marks. An indicative marking-scheme is shown in square brackets [ ] after each part of a question. A formula sheet is provided at the end of the examination paper.

**DO NOT TURN TO THE FIRST PAGE OF THE QUESTION PAPER UNTIL INSTRUCTED TO DO SO BY THE INVIGILATOR**

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1. The electric dipole contribution to the vector potential is given by

$$\mathbf{A} = -\frac{1}{4\pi\epsilon_0} \frac{e^{ikr}}{r} \frac{ik}{c} \mathbf{p}.$$

(a) Define  $k$  and  $r$  in this equation. [2 marks]

(b) What is the magnetic field  $\mathbf{B}$ ? [4 marks]

(c) Show that in the far zone  $kr \gg 1$ , with unit vector  $\mathbf{n} = \mathbf{r}/r$ ,

$$\mathbf{B} = \frac{1}{4\pi\epsilon_0} \frac{e^{ikr}}{r} \frac{k^2}{c} \mathbf{n} \times \mathbf{p},$$
$$\mathbf{E} = c \mathbf{B} \times \mathbf{n}$$

[8 marks]

(d) From the expression

$$\frac{dP}{d\Omega} = \frac{1}{2\mu_0 c} |r\mathbf{E}|^2$$

for the power radiated per unit solid angle, show that the total power radiated in all directions by pure electric dipole radiation is equal to

$$\frac{\mu_0}{12\pi c} |\mathbf{p}|^2 \omega^4$$
 [6 marks]

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2. The symmetric stress-energy-momentum tensor  $\Theta^{\alpha\beta}$  is defined by

$$\Theta^{\alpha\beta} = \frac{1}{\mu_0} [F^{\alpha\lambda} F_{\lambda}{}^{\beta} - \frac{1}{4} \eta^{\alpha\beta} F^{\mu\lambda} F_{\lambda\mu}].$$

(a) Show that

$$\Theta^{00} = \frac{1}{2} \epsilon_0 (\mathbf{E}^2 + c^2 \mathbf{B}^2),$$

and

$$\Theta^{0i} = \frac{1}{c\mu_0} (\mathbf{E} \times \mathbf{B})^i,$$

and give the physical significance of these quantities.

[8 marks]

(b) Prove

$$\partial_\alpha \Theta^{\alpha\beta} = j_\alpha F^{\alpha\beta}.$$

[6 marks]

(c) For  $\beta = 0$  show that this becomes Poynting's equation

$$\frac{\partial u}{\partial t} + \nabla \cdot \mathbf{S} = -\mathbf{E} \cdot \mathbf{j}.$$

[2 marks]

(d) What is the physical meaning of this equation?

[4 marks]

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3. (a) Consider spherical coordinates  $(r, \theta, \phi)$ , and a vector field  $\mathbf{A}$  with components

$$A_r = 0, \quad A_\theta = 0, \quad A_\phi = \frac{g(1 - \cos\theta)}{4\pi r \sin\theta}.$$

Show that this potential gives a magnetic field

$$\mathbf{B} = \frac{g}{4\pi r^2} \hat{\mathbf{r}}.$$

[5 marks]

- (b) Explain what this magnetic field describes. [3 marks]

- (c) The Lorentz force law for the motion in this field of a particle of electric charge  $e$ , rest mass  $m$ , velocity  $\mathbf{v} = \dot{\mathbf{r}} = \frac{d\mathbf{r}}{dt}$  and momentum  $\mathbf{p} = \gamma(v)m\mathbf{v}$  gives

$$\dot{\mathbf{p}} = \frac{eg}{4\pi} \frac{\mathbf{v} \times \mathbf{r}}{r^3}.$$

Show that the quantities

$$E = \sqrt{\mathbf{p}^2 c^2 + m^2 c^4}, \quad \mathbf{J} = \gamma(v)m\mathbf{r} \times \mathbf{v} - \frac{eg}{4\pi} \frac{\mathbf{r}}{r},$$

are constants of the motion and explain what these invariants are physically and what the separate terms in  $\mathbf{J}$  represent. [8 marks]

- (d) Consider the case where the particle in part (c) above is stationary. Assuming that  $\mathbf{J}$  has the properties of intrinsic angular momentum, derive the quantisation condition

$$\frac{eg}{4\pi} = \frac{n}{2}\hbar, \quad n = 0, \pm 1, \pm 2, \dots$$

[4 marks]

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4. The Liénard-Wiechert potentials for the electromagnetic fields generated by a charge  $q$  following a trajectory  $\mathbf{r} = \mathbf{r}(t)$ , with instantaneous velocity  $\mathbf{u} = \frac{d\mathbf{r}}{dt} = c\boldsymbol{\beta}$ , are

$$\Phi = \frac{q}{4\pi\epsilon_0} \left[ \frac{1}{R} \frac{1}{1 - \boldsymbol{\beta} \cdot \mathbf{n}} \right]_{\text{ret}},$$

$$\mathbf{A} = \frac{\mu_0 q c}{4\pi} \left[ \frac{\boldsymbol{\beta}}{R} \frac{1}{1 - \boldsymbol{\beta} \cdot \mathbf{n}} \right]_{\text{ret}}.$$

- (a) Explain the meaning of the notation  $[\dots]_{\text{ret}}$ , and define the distance  $R$  and the direction vector  $\mathbf{n}$ . [4 marks]

- (b) If  $|\boldsymbol{\beta}| \ll 1$ , show that at large distances from the charge the electric field is

$$\mathbf{E}_{\text{far}} = \frac{q}{4\pi\epsilon_0 c} \left[ \frac{1}{R} (\mathbf{n} \times (\mathbf{n} \times \dot{\boldsymbol{\beta}})) \right]_{\text{ret}}. \quad [6 \text{ marks}]$$

- (c) Assuming that the corresponding magnetic field is given by

$$\mathbf{B}_{\text{far}} = [\mathbf{n} \times \mathbf{E}_{\text{far}}]_{\text{far}} / c,$$

show that at large distances, the Poynting energy-flux vector is

$$\mathbf{S}_{\text{far}} = \frac{1}{\mu_0 c} |\mathbf{E}_{\text{far}}|^2 \mathbf{n}. \quad [4 \text{ marks}]$$

- (d) Derive the Larmor formula

$$P = \frac{2}{3} \frac{q^2}{4\pi\epsilon_0} \frac{1}{c^3} |\dot{\mathbf{u}}|^2$$

for the total instantaneous power radiated by a non-relativistic accelerated charge. [6 marks]

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5. (a) Show that Maxwell's equations imply that the current is conserved,  $\partial_\beta j^\beta = 0$ .  
[4 marks]

(b) Consider the gauge transformation  $A^\alpha \rightarrow A^\alpha + \partial^\alpha \Lambda$ , for some function  $\Lambda(x)$ . Explain how the Lagrangian

$$L = \int d^4x \left( -\frac{1}{4\mu_0} F^{\alpha\beta} F_{\alpha\beta} - A_\alpha j^\alpha \right)$$

is invariant under this transformation. [5 marks]

(c) Write down the equation for the Lorentz force on a particle of charge  $q$  moving with velocity  $\mathbf{v}$ . Explain how this may be generalised to the equation

$$\mathbf{f} = \rho \mathbf{E} + \mathbf{j} \times \mathbf{B}$$

for the force per unit volume on a charge and current density. [5 marks]

(d) Show that

$$f^k = F^{k\alpha} j_\alpha, \quad \text{for } k = 1, 2, 3,$$

and define  $f^0$  such that  $f^\alpha$  is a four vector. [6 marks]

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## Formula Sheet

$$\begin{aligned}
 \mathbf{a} \times (\mathbf{b} \times \mathbf{c}) &= (\mathbf{a} \cdot \mathbf{c})\mathbf{b} - (\mathbf{a} \cdot \mathbf{b})\mathbf{c}, \\
 \nabla \cdot (\psi \mathbf{a}) &= \mathbf{a} \cdot \nabla \psi + \psi \nabla \cdot \mathbf{a}, \\
 \nabla \times (\psi \mathbf{a}) &= (\nabla \psi) \times \mathbf{a} + \psi (\nabla \times \mathbf{a}), \\
 \nabla \times (\nabla \times \mathbf{a}) &= \nabla (\nabla \cdot \mathbf{a}) - \nabla^2 \mathbf{a}, \\
 \nabla (\psi(r)) &= \mathbf{n} \psi'(r).
 \end{aligned}$$

Maxwell's equations:

$$\begin{aligned}
 \nabla \cdot \mathbf{B} &= 0, & \nabla \times \mathbf{E} &= -\frac{\partial \mathbf{B}}{\partial t}, \\
 \nabla \cdot \mathbf{D} &= \rho, & \nabla \times \mathbf{H} &= \mathbf{j} + \frac{\partial \mathbf{D}}{\partial t}.
 \end{aligned}$$

$$\mathbf{F} = q(\mathbf{E} + \mathbf{v} \times \mathbf{B}).$$

For linear isotropic media:

$$\mathbf{D} = \epsilon \mathbf{E} = \epsilon_0 \mathbf{E} + \mathbf{P}, \quad \mathbf{H} = \frac{1}{\mu} \mathbf{B} = \frac{1}{\mu_0} \mathbf{B} - \mathbf{M}.$$

$$c^2 d\tau^2 = c^2 dt^2 - dx^2 - dy^2 - dz^2 = dx^\alpha \eta_{\alpha\beta} dx^\beta.$$

$$\eta_{\alpha\beta} = \begin{cases} +1 & \text{if } \alpha = \beta = 0 \\ -1 & \text{if } \alpha = \beta = 1, 2, 3 \\ 0 & \text{if } \alpha \neq \beta \end{cases}$$

$$\partial_\mu = \frac{\partial}{\partial x^\mu} = \left( \frac{1}{c} \frac{\partial}{\partial t}, \nabla \right), \quad \partial^\mu = \left( \frac{1}{c} \frac{\partial}{\partial t}, -\nabla \right).$$

$$\partial_\alpha F^{\alpha\beta} = \partial_\alpha \partial^\alpha A^\beta - \partial^\beta \partial_\alpha A^\alpha = \mu_0 j^\beta; \quad F^{\alpha\beta} = \partial^\alpha A^\beta - \partial^\beta A^\alpha.$$

$$\partial_\alpha F_{\beta\gamma} + \partial_\beta F_{\gamma\alpha} + \partial_\gamma F_{\alpha\beta} = 0.$$

$$\|F^{\alpha\beta}\| = \begin{pmatrix} 0 & -E^1/c & -E^2/c & -E^3/c \\ E^1/c & 0 & -B^3 & B^2 \\ E^2/c & B^3 & 0 & -B^1 \\ E^3/c & -B^2 & B^1 & 0 \end{pmatrix}.$$

In spherical coordinates  $(r, \theta, \phi)$ , for a vector field  $\mathbf{A}$  with components  $(A_r, A_\theta, A_\phi)$ ,

$$\begin{aligned}
 \nabla \times \mathbf{A} &= \hat{\mathbf{r}} \frac{1}{r \sin \theta} \left( \frac{\partial}{\partial \theta} (A_\phi \sin \theta) - \frac{\partial A_\theta}{\partial \phi} \right) + \hat{\theta} \left( \frac{1}{r \sin \theta} \frac{\partial A_r}{\partial \phi} - \frac{1}{r} \frac{\partial}{\partial r} (r A_\phi) \right) \\
 &\quad + \hat{\phi} \frac{1}{r} \left( \frac{\partial}{\partial r} (r A_\theta) - \frac{\partial A_r}{\partial \theta} \right)
 \end{aligned}$$