

# Answer FOUR questions

## Mark Allocation

The numbers in square brackets in the right-hand margin indicate the provisional allocation of maximum marks per sub-section of a question.

## Masses and Other Values

The following symbols may be used in this paper. The following values for these quantities may be assumed for this paper.

Meaning	Symbol	Value
Mass of $u$ quark	$m_u$	1 MeV
Mass of $d$ quark	$m_d$	2 MeV
Mass of $s$ quark	$m_s$	0.2 GeV
Mass of $c$ quark	$m_c$	1.5 GeV
Mass of $b$ quark	$m_b$	4.5 GeV
Mass of $t$ quark	$m_t$	172 GeV
Mass of all neutrinos	$m_\nu$	0
Mass of Z boson	$M_Z$	91 GeV
Mass of W boson	$M_W$	80 GeV
Width of Z boson	$\Gamma_Z$	2.5 GeV
Weinberg Angle	$\theta_w$	$28.66^\circ$
Speed of Light	$c$	$3 \times 10^8 \text{ ms}^{-1}$
Fermi Weak Decay Constant	$G_F$	$1.11 \times 10^{-5} \text{ GeV}^{-2}$
EM Coupling	$\alpha = e^2/(4\pi)$	1/137

## Dirac Matrices

The Dirac  $\gamma$  matrices satisfy  $\gamma^\mu \gamma^\nu + \gamma^\nu \gamma^\mu = 2g^{\mu\nu}$  (for  $\mu, \nu = 0,1,2,3$ ) are defined as:

$$\gamma^0 = \begin{pmatrix} I & 0 \\ 0 & -I \end{pmatrix} \quad \gamma^{i=1,2,3} = \begin{pmatrix} 0 & \sigma_i \\ -\sigma_i & 0 \end{pmatrix} \quad \gamma^5 = i\gamma^0\gamma^1\gamma^2\gamma^3 = \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix}$$

And the Pauli spin matrices,  $\sigma_i$ , are:

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

which satisfy:  $(\vec{\sigma} \cdot \vec{a})(\vec{\sigma} \cdot \vec{c}) = \vec{a} \cdot \vec{c} + i\vec{\sigma} \cdot (\vec{a} \times \vec{c})$  for 3 component vectors  $\vec{a}, \vec{c}$ .

## Cross Sections & Natural Units

$$1 \text{ barn} = 10^{-28} \text{ m}^2$$

$$\text{In natural units } 1 \text{ m} = 5.068 \times 10^{15} \text{ GeV}^{-1}.$$

1. (a) Particle A has a momentum four-vector  $P_A$  and interacts with particle B with momentum four-vector  $P_B$  to produce particles C and D with momentum four-vectors  $P_C$  and  $P_D$  respectively. The rest masses of particles A, B, C and D are  $m_A$ ,  $m_B$ ,  $m_C$  and  $m_D$  respectively. The Mandelstam variables  $s$ ,  $t$  and  $u$  are defined by:

$$s = (P_A + P_B)^2; \quad t = (P_A - P_C)^2; \quad u = (P_A - P_D)^2.$$

Show that:  $s + t + u = m_A^2 + m_B^2 + m_C^2 + m_D^2$ . [4]

- (b) Give one experimental observation for each of the three following statements that provides evidence that quarks:

- are spin 1/2 fermions [1]
- are fractionally charged [1]
- carry colour. [1]

- (c) Draw the two lowest order Feynman diagrams with the highest cross section for the scattering of a 2 GeV muon neutrino from a proton resulting in a  $\mu^-$  in the final state. [4]

- (d) The differential cross sections for scattering neutrinos,  $\sigma(\nu p)$ , and anti-neutrinos,  $\sigma(\bar{\nu} p)$ , with a stationary proton target can be approximated by:

$$\frac{d^2\sigma(\nu p)}{dxdy} = \frac{G^2 x M E}{\pi} [2d(x) + 2(1-y)^2 \bar{u}(x)]$$

$$\frac{d^2\sigma(\bar{\nu} p)}{dxdy} = \frac{G^2 x M E}{\pi} [2\bar{d}(x) + 2(1-y)^2 u(x)]$$

where  $M$  is the proton rest mass,  $E$  is the neutrino or anti-neutrino energy and  $u(x)dx$ ,  $\bar{u}(x)dx$ ,  $d(x)dx$ ,  $\bar{d}(x)dx$  represent the number of  $u$ ,  $\bar{u}$ ,  $d$  and  $\bar{d}$  quarks in the proton that carry a fractional momentum in the range  $x \rightarrow x + dx$ .

Assuming isospin symmetry between the proton and neutron and equality of their masses, show that:

$$\frac{d^2\sigma(\nu d_{np})}{dxdy} = \frac{G^2 x M E}{\pi} [u(x) + d(x) + (1-y)^2 (\bar{u}(x) + \bar{d}(x))]$$

where  $d_{np}$  is a deuteron containing one proton and one neutron and then obtain an expression for  $\frac{d^2\sigma(\bar{\nu} d_{np})}{dxdy}$ . [4]

- (e) Integrate over  $y$  and  $x$  and show that:

$$\sigma(\nu d_{np}) - \sigma(\bar{\nu} d_{np}) = \frac{2G^2 M E}{\pi}$$

[5]

2. Assume that the free particle ( $E > 0$ ) Dirac spinor solution is:

$$\psi_u^{a,b} = \sqrt{|E| + m} \begin{pmatrix} \chi^{a,b} \\ \frac{\vec{\sigma} \cdot \vec{p}}{E+m} \chi^{a,b} \end{pmatrix} e^{-ip_\mu x^\mu} \quad \chi^a = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad \chi^b = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

and the free particle ( $E < 0$ ) Dirac spinor solution is:

$$\psi_v^{a,b} = \sqrt{|E| + m} \begin{pmatrix} \frac{\vec{\sigma} \cdot \vec{p}}{E-m} \phi^{a,b} \\ \phi^{a,b} \end{pmatrix} e^{-ip_\mu x^\mu} \quad \phi^a = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \quad \phi^b = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

(a) The charge conjugate spinor of  $\psi$ ,  $\psi_C$ , is defined by

$$\psi_C = i\gamma^2 \psi^* .$$

By using the explicit form of the Pauli matrices, determine the charge conjugate spinor of  $\psi_u^a$  and show how it is related to  $\psi_v^a$ . Explain the significance of this result.

[8]

(b) Show that if  $m \ll E$

$$\gamma^5 \psi_u^a \sim \begin{pmatrix} \vec{\sigma} \cdot \hat{p} & 0 \\ 0 & \vec{\sigma} \cdot \hat{p} \end{pmatrix} \psi_u^a$$

where  $\hat{p} = \vec{p}/|\vec{p}|$

[4]

(c) If the projection operators,  $P_R$  and  $P_L$ , are defined by:

$$P_R = \frac{1}{2}(1 + \gamma^5) \quad P_L = \frac{1}{2}(1 - \gamma^5)$$

Show, by constructing a helicity operator, that in the ultra-relativistic limit,  $P_{R,L}$ , project out the positive and negative helicity components of a free Dirac spinor.

[5]

(d) What experimental observation could be used to show that the neutrino is a Majorana and not a Dirac particle ?

Explain why the observation is not allowed for Dirac particles.

[3]

3. (a) What measurements at the Tevatron  $p\bar{p}$  collider can be used to constrain the predicted mass of the Higgs boson? [2]
- (b) Draw the Feynman diagram that has the highest cross section for the production and subsequent decay of a Higgs boson of mass 160 GeV at the Tevatron collider. You need not consider any hadronisation processes. [5]
- (c) Draw a Feynman diagram for a process that will occur at a far greater rate than the above Higgs process but will result in the same final state particles. Explain briefly why the rate is so much higher. [4]
- (d) If the Higgs boson decays to two particles  $A$  and  $B$  of mass  $m_A$  and  $m_B$  respectively, show that the invariant mass,  $m_{\text{inv}}$ , of the  $A + B$  system and hence of the Higgs boson is given by:

$$m_{\text{inv}}^2 = m_A^2 + m_B^2 + 2 \left[ E_T^A E_T^B \cosh(\eta_A - \eta_B) - \mathbf{p}_T^A \cdot \mathbf{p}_T^B \right]$$

where  $E_T \equiv \sqrt{E^2 - p_z^2}$ ,  $\mathbf{p}_T$  is the transverse momentum 2-vector =  $(p_x, p_y)$ .

You may assume that:

$$E_T \cosh \eta = E, \quad E_T \sinh \eta = p_z, \quad \cosh(x + y) = \cosh x \cosh y + \sinh x \sinh y$$

[4]

- (e) If the angle between  $\mathbf{p}_T^A$  and  $\mathbf{p}_T^B$  is  $\Delta\phi$  and  $\Delta\eta = \eta_A - \eta_B$ , show, using the appropriate Taylor expansions, that for massless  $A$  and  $B$  particles with small  $\Delta\phi$  and  $\Delta\eta$

$$m_{\text{inv}}^2 \approx |\mathbf{p}_T^A| |\mathbf{p}_T^B| (\Delta\eta^2 + \Delta\phi^2)$$

You may assume that:

$$\cos(x) = \sum_{n=0}^{\infty} (-1)^n \frac{x^{2n}}{(2n)!}; \quad \cosh(x) = \cos(ix)$$

[3]

- (f) If no Higgs boson is found at the LHC; draw a Feynman diagram of the process one could study to try to elucidate the mechanism of electroweak symmetry breaking. [2]

4. (a) Explain with the help of two Feynman diagrams why the value of the Fermi weak decay constant,  $G_F^\beta$ , deduced from the rate of nuclear  $\beta^-$  decays is slightly less than the value,  $G_F^\mu$ , deduced from the rate of  $\mu^-$  decay. Approximately what value would you expect for  $G_F^\beta/G_F^\mu$ ? [4]
- (b) Explain why the introduction of a phase into the CKM matrix can produce CP violation and why observations of CP violation are important. [3]
- (c) Write down the formula for the partial width,  $\Gamma_{cb}$ , for the decay  $\overline{B}^0 \rightarrow D^+ \mu^- \overline{\nu}_\mu$  in terms of the  $\overline{B}^0$  lifetime,  $\tau_B$ , and the branching ratio,  $BR$ , for the decay. How is  $\Gamma_{cb}$  related to  $V_{cb}$ ? [2]
- (d) Draw a Feynman diagram for the decay  $\overline{B}^0 \rightarrow D^+ \mu^- \overline{\nu}_\mu$ , followed by the decay  $D^+ \rightarrow \overline{K}_s^0 \pi^+$ . Explain briefly, with reference to appropriate particle detectors, how one could identify the  $D^+$  meson in this decay sequence. [5]
- (e) The LEP accelerator had 2 counter circling beams that collided electrons and positrons head on. Both beams had an energy of 45.5 GeV. Consider the case of  $Z$  production and its subsequent decay to a  $b\overline{b}$  pair. Assuming the  $b$  quarks form B-meson bound states, how far on average (in cm) would one expect each of the B-mesons to travel before decay assuming that the lifetime of B mesons is  $\sim 1.5$  ps? How could one measure this decay distance? [6]

The quark content of  $\overline{B}^0$  is  $b\overline{d}$ ,  $D^+$  is  $c\overline{d}$ ,  $\overline{K}_s^0$  is  $s\overline{d}$  and  $\pi^+$  is  $u\overline{d}$

5. (a) Draw the two lowest order Feynman diagrams for  $e^+e^- \rightarrow \mu^+\mu^-$ .  
Write down fermion coupling expressions for the vertex factors in each diagram in terms of the electromagnetic coupling,  $g$ , the Dirac gamma matrices, the Weinberg angle,  $\theta_W$ , and the vector,  $C_{fV}$ , and the axial-vector,  $C_{fA}$  couplings. [4]
- (b) If  $\theta$  is the angle of the  $\mu^-$  with respect to the incoming  $e^-$  in the  $e^+e^-$  centre of mass frame, how would one expect the cross section for the above process to depend on  $\cos\theta$  for a purely vector interaction? Why is this not observed even for  $e^+e^-$  centre of mass energies,  $\sqrt{s}$ , of 30 GeV? [3]
- (c)  $A_{FB}$  is a measurement of the angular asymmetry in  $\cos\theta$  of  $\mu^-$  from the above interaction. With reference to the Fermi weak decay constant,  $G_F$ , the EM coupling,  $\alpha$ , and the centre of mass energy,  $\sqrt{s}$ , obtain a simple expression for the approximate value of  $A_{FB}$  for  $\sqrt{s} \ll M_z$  and determine a value of  $A_{FB}$  at  $s = 900 \text{ GeV}^2$ . [5]
- (d) Why do we expect the above formulae for  $A_{FB}$  not to be valid for the process  $e^+e^- \rightarrow e^+e^-$  at  $\sqrt{s} = 30 \text{ GeV}$ ? [2]
- (e) The above formulae for  $A_{FB}$  are leading order formulae. Draw a higher order Feynman diagram that would modify the prediction for  $A_{FB}$  and could be used to make a prediction for the mass of the top quark. [2]
- (f) At the Tevatron  $p\bar{p}$  collider, W bosons are produced through quark anti-quark annihilation.  $\theta$  is the angle defined with respect to the proton direction. On average, the  $u$  quark carries a greater fraction of the proton's momentum than the  $d$  quark. With reference to appropriate Feynman diagrams, explain what one would expect for the distribution of positrons and electrons from W decay as a function of  $\cos\theta$ . [4]

6. (a) The Lagrangian density for free electrons is:

$$\mathcal{L} = i\bar{\psi}\gamma^\mu\partial_\mu\psi - m\bar{\psi}\psi$$

Show that  $\mathcal{L}$  is invariant under global gauge (phase) transformations but not under local gauge (phase) transformations. [6]

- (b) Local gauge invariance can be achieved by replacing the derivative  $\partial_\mu$  with the “covariant derivative”

$$\mathcal{D}_\mu \equiv \partial_\mu - ieA_\mu$$

( $e$  is a constant), as long as  $A_\mu$  transforms in a certain way. Derive the transformation rules for  $A_\mu$  and give its physical interpretation. [6]

- (c) What would be the implications of adding a term of the form  $\frac{1}{2}m_A^2 A^\mu A_\mu$  to the QED Lagrangian density? [3]

- (d) Show that the substitution of the Lagrangian density:

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - j_\mu A^\mu$$

into the Euler-Lagrange equation for  $A_\mu$  gives Maxwell’s equations i.e.

$$\partial_\mu F^{\mu\nu} = j^\nu$$

where  $F^{\mu\nu} = \partial^\mu A^\nu - \partial^\nu A^\mu$ . [5]

The Euler-Lagrange equation for a variable  $\phi$  is:

$$\partial_\mu \left( \frac{\partial \mathcal{L}}{\partial(\partial_\mu \phi)} \right) - \frac{\partial \mathcal{L}}{\partial \phi} = 0$$